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SUSY Searches at the Tevatron

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Abstract

In this article results from supersymmetry searches at D0 and CDF are reported. Searches for third generation squarks, searches for gauginos, and searches for models with R-parity violation are described. As no signs of supersymmetry for these models are observed, the most stringent limits to date are presented.

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1 Introduction

The standard model (SM) describes matter and its interactions at the elementary level. However it is not a complete theory, since there are indications of phenomena which the SM does not describe. Supersymmetry (SUSY) [1] is one of the extensions of the SM that provides a solution to many shortcomings of the SM. SUSY is the symmetry between different states of spin, i.e. between fermions and bosons. Since each SM particle has its SUSY partner, sparticle, the number of parameters in SUSY is increased. The minimal supersymmetric standard model (MSSM) is a model with minimal particle content. It is possible in the MSSM that the third generation squarks are lighter than the first and second, thus making them accessible to Tevatron experiments.

It is known that particles and sparticles do not have the same mass, which is an indication that SUSY is broken. There are several possible mechanisms of SUSY breaking. D0 and CDF have searched for gauginos within the minimal supergravity model (mSUGRA), and gauge mediated symmetry breaking model (GMSB). The mSUGRA is described by five independent parameters: the unified scalar and gaugino masses m_0 and $m_{1/2}$, the ratio of the vacuum expectation values of the two Higgs doublets $\tan\beta$, the unified trilinear coupling A_0 , and the sign of the Higgs mass parameter $\text{sign}(\mu)$. GMSB is described by six parameters: the breaking scale Λ , the mass and number of messengers M and N , the scale factor of gravitino mass C_{grav} , $\tan\beta$ and the $\text{sign}(\mu)$.

R-parity, defined as $R = (-1)^{3(B-L)+2S}$, where B is baryon number, L is lepton number, and S is spin, is introduced to preserve lepton and baryon number conservation. If it is conserved the lightest supersymmetric particle (LSP) is stable, all sparticles eventually decay to LSP, and sparticles are produced in pairs. However, lepton and baryon number symmetry is accidental, so R-parity is imposed by hand. If R-parity is violated SUSY particles decay to SM particles and they can be produced singly, which leads to higher cross sections at Tevatron.

The CDF [2] and D0 [3, 4, 5] detectors are multi purpose detectors. They consist of central tracking system, electromagnetic and hadronic calorimeters and outer muon detectors. They recorded about 10 fb^{-1} of data between April 2002 and May 2011. Results presented in this article use up to 6.3 fb^{-1} .

2 Third generation squark searches

The D0 experiment searched for pair production of sbottom quarks [6] in 5.2 fb^{-1} of data. It was assumed that the mass of the sbottom quark satisfied $m_b + m_{\tilde{\chi}_1^0} < m_{\tilde{b}_1} < m_t + m_{\tilde{\chi}_1^-}$, so only $\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ was possible, and that $\tilde{\chi}_1^0$ was the LSP. The search was interpreted in the MSSM with R-parity conservation. The signature of this final

state was two b jets and large missing transverse energy \cancel{E}_T , so data were recorded using a combination of jet and \cancel{E}_T based triggers. The dominant SM background was from multijet (MJ) events where instrumental \cancel{E}_T was present due to jet resolution and mismeasurement. It was estimated from data. Other backgrounds included W/Z +jets, $t\bar{t}$ and diboson events, which were obtained from Monte Carlo (MC) simulation. Events were required to have at least two not back-to-back b -tagged jets with $p_T > 20$ GeV and $\cancel{E}_T > 40$ GeV. Selection was further optimized for different choices of $(m_{\tilde{b}_1}, m_{\tilde{\chi}_1^0})$ to achieve the maximal sensitivity. Good agreement between data and SM backgrounds was observed as can be seen in Fig. 1 (left) and Table 1. A 95% C.L limit on the production of sbottom quark pairs was set and the excluded region is shown in the plane of the bottom squark versus neutralino mass in Fig. 1 (right).

$m_{\tilde{b}_1}, m_{\tilde{\chi}_1^0}$	(130,85)	(240,0)
Data	901	7
Background	971 ± 52	6.9 ± 1.7
Signal	481 ± 66	10.5 ± 1.9

Table 1: Number of events for the two chosen $(m_{\tilde{b}_1}, m_{\tilde{\chi}_1^0})$ points for data, SM expectation and sbottom quark signal.

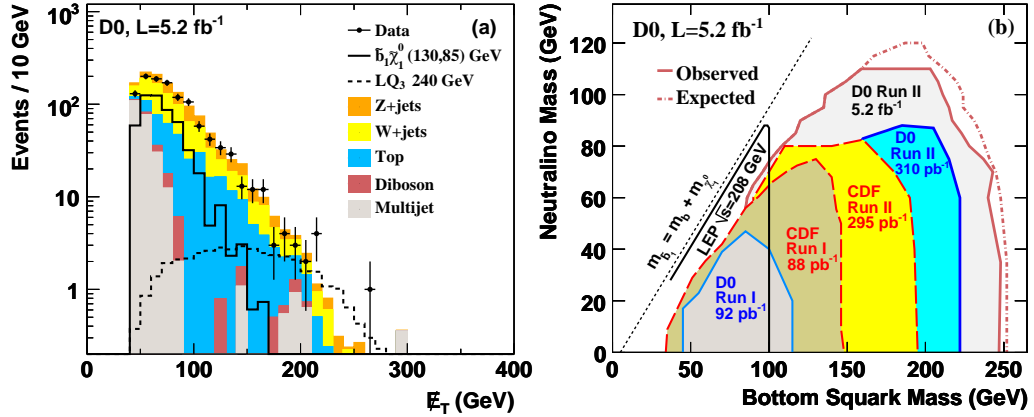


Figure 1: The \cancel{E}_T distribution in the $\tilde{b}_1\tilde{b}_1 \rightarrow b\tilde{\chi}_1^0\bar{b}\tilde{\chi}_1^0$ search for data, expected SM backgrounds, and sbottom quark signal for $m_{\tilde{b}_1} = 130$ GeV and $m_{\tilde{\chi}_1^0} = 85$ GeV (left). The observed and expected 95 % C.L. exclusion contour in a $m_{\tilde{b}_1}$ and $m_{\tilde{\chi}_1^0}$ plane (right).

D0 [7] and CDF [8] searched for pair production of scalar top quarks. D0 searched

for stop quark pair production $\tilde{t}_1\bar{\tilde{t}}_1$ with the $b\bar{b}e^\pm\mu^\mp\tilde{\nu}\bar{\tilde{\nu}}$ final state in 5.4 fb^{-1} of data. It was assumed that the branching ratio for $\tilde{t}_1 \rightarrow b\tilde{\nu}$ was 100% and that the sneutrino $\tilde{\nu}$ was either the LSP or that it decayed invisibly into a neutrino and a neutralino. Events were selected with one isolated electron with $p_T > 15 \text{ GeV}$ and $|\eta| < 1.1$ [9], one isolated muon with $p_T > 10 \text{ GeV}$ and $|\eta| < 2$, and $\cancel{E}_T > 7 \text{ GeV}$ at the preselection level. The dominant SM backgrounds included $Z \rightarrow \tau\tau$ where both τ leptons decayed leptonically, top quark pairs, diboson production, W +jets and instrumental MJ background. To further optimize this search, samples were divided according to the mass difference between \tilde{t}_1 and $\tilde{\nu}$, $\Delta M = M_{\tilde{t}_1} - M_{\tilde{\nu}}$. D0 chose two benchmark points, $(M_{\tilde{t}_1}, M_{\tilde{\nu}}) = (200, 100) \text{ GeV}$ for "large- ΔM ", $\Delta M > 60 \text{ GeV}$, and $(M_{\tilde{t}_1}, M_{\tilde{\nu}}) = (110, 90) \text{ GeV}$ for "small- ΔM ", $\Delta M < 60 \text{ GeV}$. Signal selection was optimized as a function of ΔM . In a dominant background after preselection, $Z \rightarrow \tau\tau$, two leptons were often back to back, and \cancel{E}_T was low. Thus, an additional requirement was introduced, events were rejected if $\Delta\phi(e, \mu) > 2.8$ and $\cancel{E}_T < 20 \text{ GeV}$. Fig. 2 (left) shows muon p_T after this selection. To discriminate against dominant backgrounds ($Z \rightarrow \tau\tau$, $t\bar{t}$ and diboson) and signal, a separate multivariate analysis output (MVA) was built for each background. A cut was then applied on the most discriminating MVA, which was $t\bar{t}$ for "small- ΔM ", and $Z \rightarrow \tau\tau$ for "large- ΔM ". A two dimensional histogram of the remaining two MVA distributions was used to search for signal. In the absence of any significant excess, the 95% C.L. exclusion limits on scalar top quark production as a function of the $\tilde{\nu}$ and \tilde{t}_1 masses was set, and is shown in Fig. 2 (right).

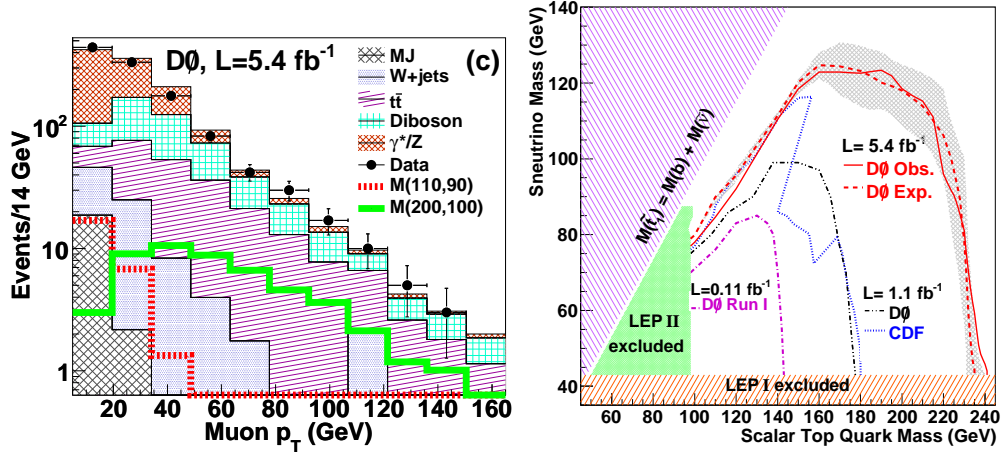


Figure 2: The p_T of the muon for data, expected background and two different $M_{\tilde{t}_1}$ representing small- ΔM and large- ΔM (left). The observed and expected 95% C.L. exclusion contours in the $(m_{\tilde{t}_1}, m_{\tilde{\nu}})$ plane (right).

CDF searched for the pair production of scalar top quarks further decaying to

the final state with two leptons, $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm \rightarrow b\tilde{\chi}_1^0 l\nu$ in 2.7 fb^{-1} of data. It was assumed $m_{\tilde{t}_1} < m_t$, neutralino $\tilde{\chi}_1^0$ is LSP, $m_{\tilde{\chi}_1^\pm} < m_{\tilde{t}_1} - m_b$ and R-parity conservation. Depending on the masses of chargino, $\tilde{\chi}_1^\pm$, and neutralino $\tilde{\chi}_1^0$, and chosen SUSY parameters, the branching ratio $BR(\tilde{\chi}_1^\pm \rightarrow b\tilde{\chi}_1^0 l\nu)$ can have values between 0.11 and 1.0. The final state, containing two leptons, two b quarks and \cancel{E}_T from neutralino is identical to top quark pair production. Thus the reconstructed top mass will be biased toward lower values in the presence of scalar top quarks. Events were selected with two leptons with $p_T > 20 \text{ GeV}$, two jets with $p_T > 12 \text{ GeV}$, and $\cancel{E}_T > 20 \text{ GeV}$. The dominant SM backgrounds were $t\bar{t}$, $Z + \text{jets}$, diboson and instrumental MJ background. To maximize sensitivity of this search, samples were divided into events with no- b -tagged jets and at least one b -tagged jet. Data was in agreement with SM prediction as shown in Table 2. Scalar top yield was calculated for $m_{\tilde{t}_1} = 132.5 \text{ GeV}$, $m_{\tilde{\chi}_1^\pm} = 105.8 \text{ GeV}$, $m_{\tilde{\chi}_1^0} = 47.6 \text{ GeV}$ and $BR = 0.11$. Figure 3 (left) shows the reconstructed stop quark mass for $m_{\tilde{t}_1} = 138.3 \text{ GeV}$, $m_{\tilde{\chi}_1^\pm} = 105.8 \text{ GeV}$, $m_{\tilde{\chi}_1^0} = 76.2 \text{ GeV}$ and $BR = 0.5$. In the absence of a significant excess in data mass limits in the $(m_{\tilde{\chi}_1^0}, m_{\tilde{t}_1})$ plane for different branching ratios was set (see Fig. 3 (right)).

	Data	Total MC	Signal
No b -tag	65	65.9 ± 9.8	3.9 ± 0.9
$\geq 1b$ -tag	57	56.4 ± 7.2	9.5 ± 1.9

Table 2: Event yields for data, predicted backgrounds and expected signal with no b -tagged jets and with at least one b -tagged jet.

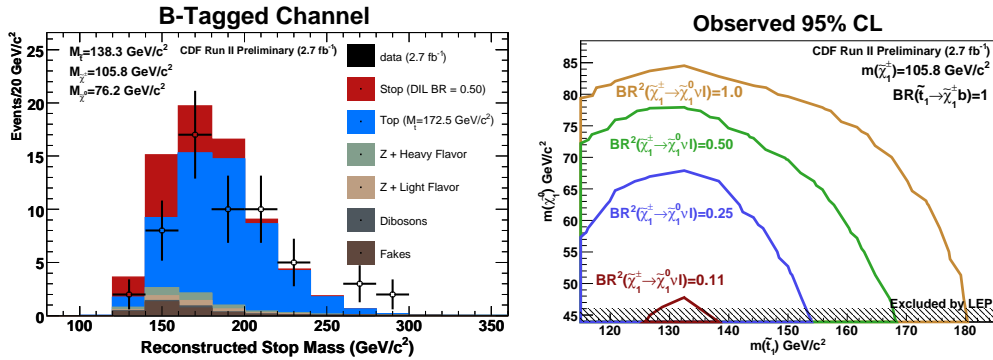


Figure 3: Reconstructed $m_{\tilde{t}_1}$ for data, predicted background and expected signal for events with at least one b -tagged jet (left). The observed 95% C.L. limits in $(m_{\tilde{t}_1}, m_{\tilde{\chi}_1^0})$ plane for various branching ratios to the dilepton final state (right).

3 Gaugino Searches

CDF [10] and D0 [11] searched for the production of charginos and neutralinos in the final states with at least three leptons in 3.2 fb^{-1} (CDF) and 2.3 fb^{-1} (D0) of data. At CDF events were divided in search channels based on a lepton flavor and quality. Electrons and muons can be tight or loose, and tracks were always tight. Each event was categorized into an exclusive trilepton channel composed of combinations of these objects. All objects were required to be central, $|\eta| < 1$, and isolated from other objects. Opposite sign leptons were required not to be back to back, and their invariant mass to be outside the Z mass window. In addition events were selected if they have less than two jets. The invariant mass after all cuts except the cut on invariant mass itself is shown in Fig. 4 (left). Results were interpreted within mSUGRA. The exclusion contour in a $(m_0, m_{1/2})$ plane is shown in Fig. 4 (right).

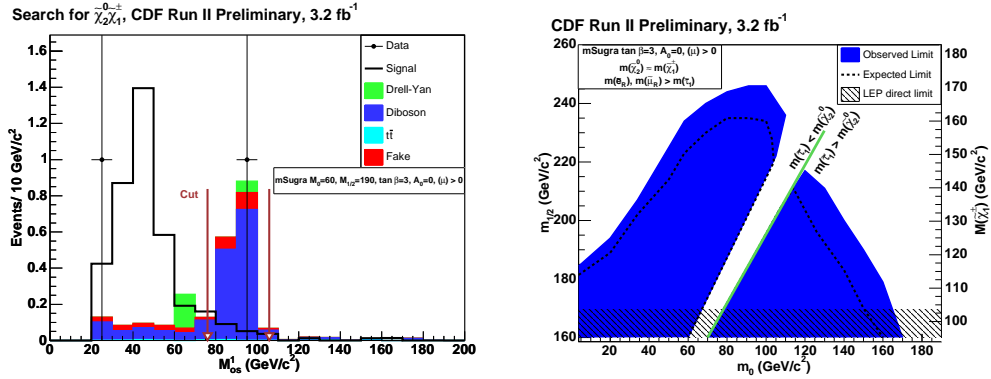


Figure 4: Invariant mass of the two opposite sign leptons after preselection cuts (left). Expected and observed 95% C.L. limit contours for mSUGRA in a $(m_0, m_{1/2})$ plane (right).

D0 defined four different channels distinguished by the lepton content of the final state, dielectron plus lepton (eel), dimuon plus lepton ($\mu\mu l$), electron, muon plus lepton ($e\mu l$), and muon, τ lepton, where τ lepton was identified through its hadronic decays, plus lepton. For each channel, one "low- p_T " and one "high- p_T " selection was designed to maximize sensitivity for various parameter points in the $(m_0, m_{1/2})$ plane. Selection requirements included $p_T(l_1) > 12 \text{ GeV}$, $p_T(l_2) > 8 \text{ GeV}$, rejecting events with same flavor dilepton mass in a Z mass window, rejecting events with back-to-back leptons, and requiring large \cancel{E}_T . After all selection requirements good agreement between data and SM backgrounds was observed as shown for the invariant mass of two electrons in Fig. 5 (left). Limits on chargino and neutralino production were set and interpreted in mSUGRA. The excluded region in the $(m_0, m_{1/2})$ plane is shown in Fig. 5 (right).

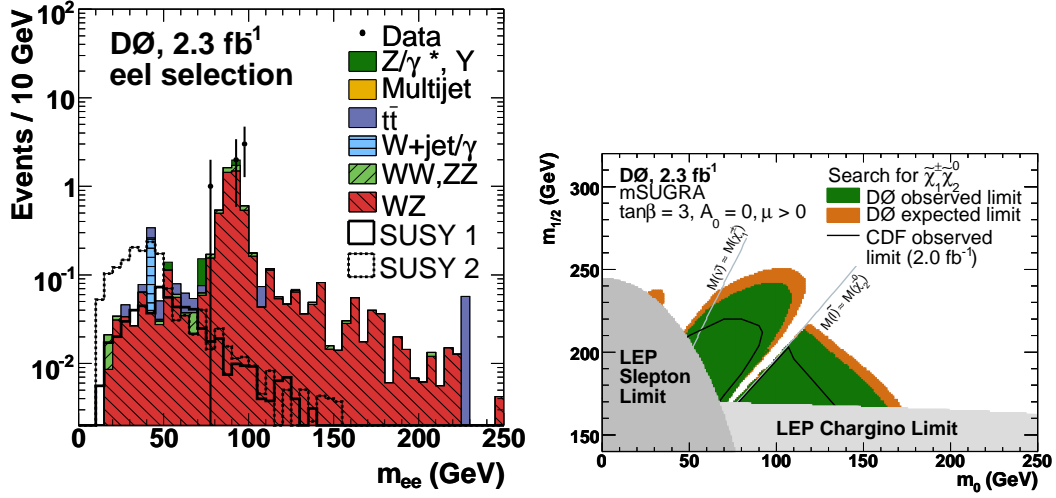


Figure 5: Invariant mass of the two electrons after all cuts except the one on m_{ee} for data, expected background and SUSY signal (left). The expected and observed 95% C.L. limit contours in the $(m_0, m_{1/2})$ plane (right).

CDF searched for SUSY signatures in events with two leptons with the same electric charge [12] in 6.1 fb $^{-1}$ of data. A simplified SUSY model was built, where only particles that appear in this model have been tested. This model was based on the following requirements:

- For final states with leptons, the model contains chargino $\tilde{\chi}_1^\pm$ and neutralino $\tilde{\chi}_1^0$ which decay to W and Z bosons. Slepton modes were not considered.
- R-parity conservation and the presence of a stable LSP were assumed.
- The largest cross section will be for pair production of colored states squarks and gluinos.

Events were selected with at least two same sign leptons with $p_T > 15$ GeV for the leading and $p_T > 10$ GeV for second leading lepton, and at least two jets with $p_T > 15$ GeV. The scalar sum of the leptons and jets p_T (H_T) is shown in Fig. 6 (left). No significant excess in data was observed and limits on gluino and squark production were set. Figure 6 (right) shows limit on squark pair production for LSP mass of 100 GeV. Excluded area is above the curve.

DØ searched for SUSY in GMSB models where gravitino \tilde{G} is very light and thus is the LSP and the lightest neutralino $\tilde{\chi}_1^0$ is the NLSP [13] in 6.3 fb $^{-1}$ of data. The largest cross section at Tevatron, assuming R-parity conservation, is chargino and neutralino pair production ($\tilde{\chi}_1^\pm \tilde{\chi}_2^0, \tilde{\chi}_1^\pm \tilde{\chi}_1^\pm$) with subsequent decay chains to NLSP $\tilde{\chi}_1^0$. A case where $\tilde{\chi}_1^0$ decays promptly to photon and an essentially massless gravitino,

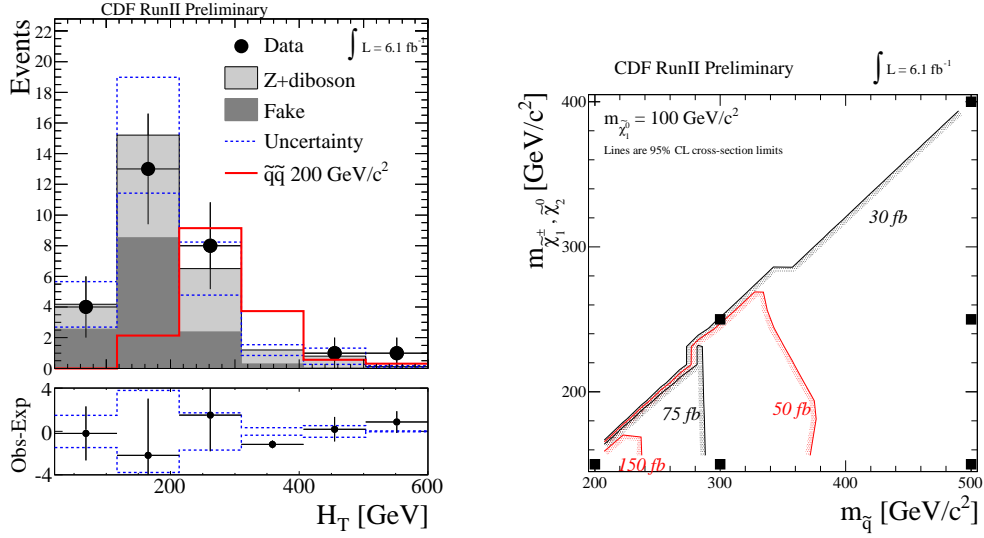


Figure 6: The scalar sum of the leptons and jets p_T , H_T , in the events with the two same sign leptons and at least two jets (left). 95% C.L. limits on squark pair production for LSP mass of 100 GeV (right). Excluded is the area above the curve.

$\tilde{\chi}_1^0 \rightarrow \tilde{G}\gamma$, was considered. Events were selected with at least two isolated photons in central calorimeter region with $p_T > 25$ GeV. To remove events with poorly modeled \cancel{E}_T , events were rejected with $\Delta\phi$ between \cancel{E}_T and leading jet (if present) greater than 2.5, and $\Delta\phi$ between \cancel{E}_T and any photon less than 0.2. SM backgrounds, which were categorized as instrumental \cancel{E}_T sources ($\gamma\gamma$, γ +jet and MJ), and genuine \cancel{E}_T sources ($W\gamma$, W +jets, $W/Z + \gamma\gamma$), were estimated from data. Figure 7 (left) shows the \cancel{E}_T distribution in $\gamma\gamma$ events, where a good agreement between data and SM backgrounds was observed. A GMSB scenario was probed using the set of parameters from the SPS8 model, where the scale Λ was unconstrained, $M_{mes} = 2\Lambda$, $N_{mes} = 1$, $\tan\beta = 15$, and $\mu > 0$. Since no evidence of the SUSY in the \cancel{E}_T distribution was observed, a 95% C.L. upper limit on the cross section production as a function of Λ , $m_{\tilde{\chi}_1^0}$ and $m_{\tilde{\chi}_1^\pm}$ was set (Fig. 7 (right)).

D0 searched for signatures of hidden-valley models [14] in 5.8 fb^{-1} of data. Hidden valley models contain a hidden sector that is very weakly coupled to the SM. The force carrier in the hidden-valley model is the dark photon (γ_D), which is assumed to be very light ($m_{\gamma_D} < 2$ GeV) and decays into a pair of SM charged fermions. In these models, SUSY will have partners for both the SM and the hidden sector. It was assumed that the LSP of the hidden sector \tilde{X} was lighter than the lightest SM SUSY partner SM-LSP. Then the SM-LSP decayed promptly into particles of the hidden sector if R-parity was conserved. Hidden-sector particles are light so their

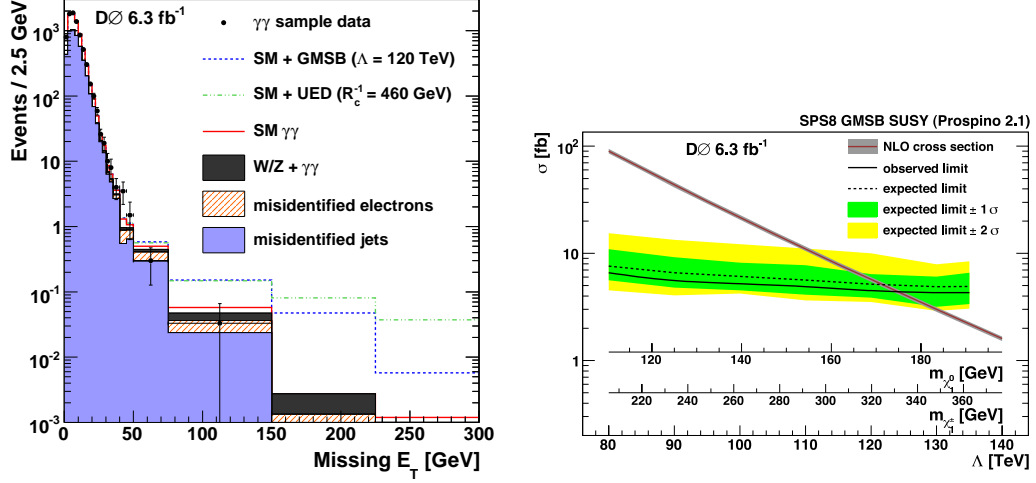


Figure 7: \cancel{E}_T distribution in $\gamma\gamma$ events for data, predicted backgrounds and expected signal (left). The expected and observed 95% C.L. upper cross section limit for GMSB as a function of Λ , $m_{\tilde{\chi}_1^0}$ and $m_{\tilde{\chi}_1^\pm}$ (right).

decays produce jets of tightly collimated particles from decays of γ_D , i.e. jets of leptons or leptonic jets (l -jets). Since γ_D can decay to a pair of electrons or to a pair of muons, l -jet can be either "electron l -jet" (e_{lj}) or "muon l -jet" (μ_{lj}). Events were selected with two l -jets, $e_{lj}e_{lj}$, $e_{lj}\mu_{lj}$ or $\mu_{lj}\mu_{lj}$. The main backgrounds, which were estimated from a data, consist of MJ production, and, in a case of e_{lj} , photon production with a decay to e^+e^- pairs. We searched for presence of leptonic jets in the events with $\cancel{E}_T > 30$ GeV. Figure 8 (left) shows invariant mass of two muon l jets for different assumption of m_{γ_D} . With no significant excess observed in data, 95% C.L. upper limits were set as a function of m_{γ_D} on the production cross section for SUSY particles decaying to two l -jets and large \cancel{E}_T (Fig. 8 (right)).

4 R-parity violation searches

If R-parity is violated single production of SUSY particles is allowed, and they can be observed at Tevatron. D0 searched for resonant production of a sneutrino which decays into an electron and a muon in 5.4 fb^{-1} of data [15]. The relevant terms in Lagrangian which describe production and decays of $\tilde{\nu}$ are:

$$\mathcal{L}_{RPV} = -\frac{1}{2}\lambda_{ijk}(\tilde{\nu}_{iL}\bar{l}_{kR}l_{jL} - \tilde{\nu}_{jL}\bar{l}_{kR}l_{iL}) - \lambda'_{ijk}(\tilde{\nu}_{iL}\bar{d}_{kR}d_{jL}) + h.c. \quad (1)$$

where $i, j, k = 1, 2, 3$ is fermion generation, l is down type lepton field and d is down type quark field. Thus at Tevatron, single $\tilde{\nu}$ can be produced in $d\bar{d}$ scattering. It

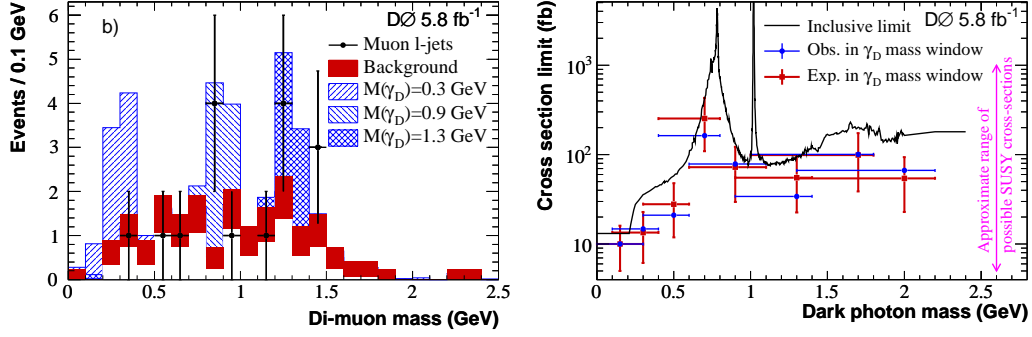


Figure 8: Invariant mass of dark photon candidates with two isolated l -jets and $\cancel{E}_T > 30$ GeV where γ_D decayed to two muons (left). 95% C.L. limits on the cross section as a function of the m_{γ_D} (right).

was assumed that only the third generation sneutrino ($\tilde{\nu}_\tau$) was produced, and that it was the LSP. Further it is assumed that all the couplings are equal to zero except λ'_{311} and $\lambda_{312} = \lambda_{321} = -\lambda_{231} = -\lambda_{132}$. The dominant background processes for this search are $Z \rightarrow \tau\tau$, $t\bar{t}$, diboson and $W + jets$. Events were selected with one isolated electron with $p_T > 30$ GeV, one isolated muon with $p_T > 25$ GeV and no jets with $p_T > 25$ GeV. Signal events have low \cancel{E}_T , but due to limited momentum resolution of the muons, some \cancel{E}_T will be present if it is either aligned or anti-aligned with muon. Thus events were rejected with $\cancel{E}_T > 20$ GeV and $0.7 < \Delta\phi(\mu, \cancel{E}_T) < 2.3$. Figure 9 (left) shows $M_{e\mu}$ where good agreement between data and SM background was observed. With no significant excess present in a data 95% C.L. upper limits were set on couplings as a function of $M_{\tilde{\nu}_\tau}$, as shown in Fig. 9 (right).

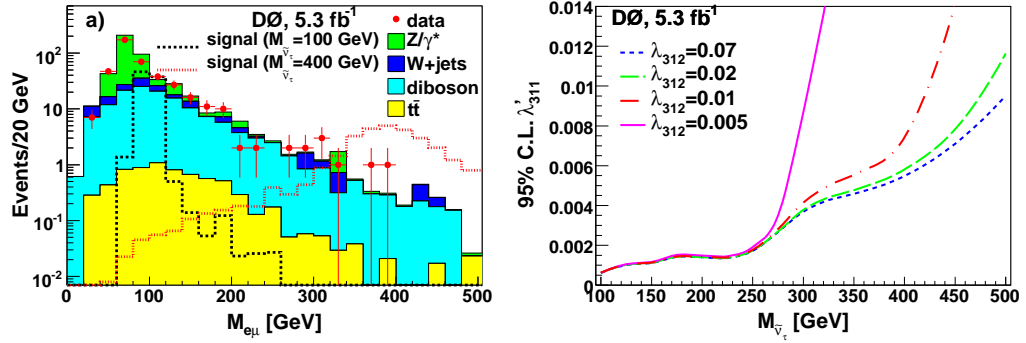


Figure 9: Distribution of $M_{e\mu}$ for data, predicted background and expected signal (left). The 95% C.L. observed upper limits on λ'_{311} for values of λ_{312} as a function of $M_{\tilde{\nu}_\tau}$ (right).

CDF [16] searched for 3-jets resonances in 3.2 fb^{-1} of data. Although this search

was model independent, the possible new physics was modeled with RPV SUSY gluino pair production, each gluino decaying into three partons. The dominant background was MJ and it was estimated from data. Events were selected with $\cancel{E}_T < 50$ GeV, between one and four primary vertices, and at least six jets. It was required that the sum of the p_T of the six leading jets was greater than 250 GeV. The final requirement $\sum_{jjj} p_T - M_{jjj} > \text{offset}$, with distributions shown in Fig. 10 (left) for RPV gluino with mass of 190 GeV and Fig. 10 (right) for data, was optimized for different mass points. MJ background was estimated from five jets data and fitted with Landau function. To extract signal from combinatorial background Landau+Gaussian fit was used (see Fig. 11 (left)), and number of signal events was obtained by integrating the Gaussian in $\pm 1\sigma$ range. No significant excess was observed in data, and a 95% C.L. limit on the cross section times branching ratio as a function of gluino mass was set, and it is shown in Fig. 11 (right).

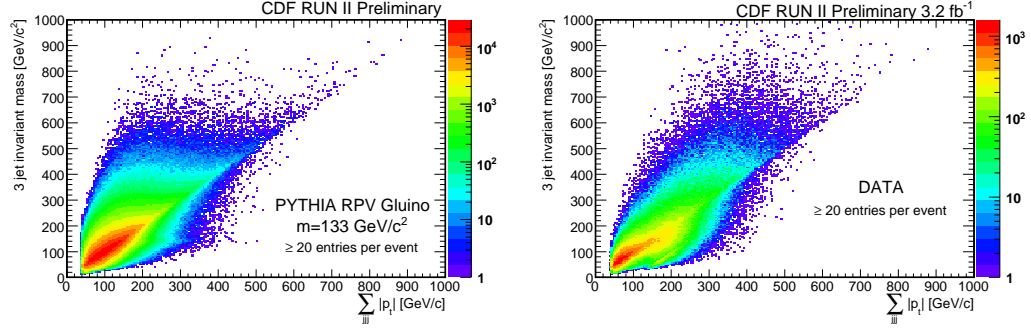


Figure 10: Distributions of the M_{jjj} vs $\sum_{jjj} p_T$ for the RPV gluino signal (left) and data (right). For each events there are multiple entries (> 20).

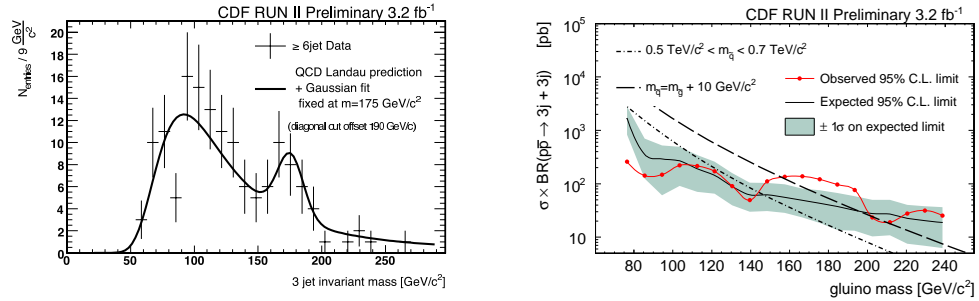


Figure 11: Final mass distribution for the data and QCD Landau prediction plus signal Gaussian fit (left). The 95% C.L. expected and observed limit (right).

5 Conclusion

Results from SUSY searches at the Tevatron have been presented. It is shown that data agrees with SM expectations, and thus SUSY parameter space has been constrained. Limits set by various CDF and D0 searches are the most stringent to date.

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References

- [1] S. P. Martin, arXiv:hep-ph/9709356.
- [2] A. Abulencia *et al.*, CDF Collaboration, J. Phys. G: Nucl. Part. Phys. **34**, (2007).
- [3] V. M. Abazov *et al.*, D0 Collaboration, Nucl. Instrum. Methods Phys. Res. A **565**, 463 (2006).
- [4] M. Abolins *et al.*, D0 Collaboration, Nucl. Instrum. Methods Phys. Res. A **584**, 75 (2008).
- [5] R. Angstadt *et al.*, D0 Collaboration, Nucl. Instrum. Methods Phys. Res. A **622**, 298 (2010).
- [6] V. M. Abazov *et al.*, D0 Collaboration, Phys. Lett. B **693**, 92 (2010).

- [7] V. M. Abazov *et al.*, D0 Collaboration, Phys. Lett. B **696**, 321 (2011).
- [8] T. Aaltonen *et al.*, CDF Collaboration, Phys. Rev. Lett. **104**, 251801 (2010).
- [9] The pseudorapidity is defined as $\eta = -\ln[\tan(\theta/2)]$, where θ is the polar angle with respect to the proton beam direction.
- [10] T. Aaltonen *et al.*, CDF Collaboration, CDF Public Note 9817.
- [11] V. M. Abazov *et al.*, D0 Collaboration, Phys. Lett. B **680**, 34 (2009).
- [12] T. Aaltonen *et al.*, CDF Collaboration, CDF/PHYS/EXO/PUBLIC/10464.
- [13] V. M. Abazov *et al.*, D0 Collaboration, Phys. Rev. Lett. **105**, 221802 (2010).
- [14] V. M. Abazov *et al.*, D0 Collaboration, Phys. Rev. Lett. **105**, 211802 (2010).
- [15] V. M. Abazov *et al.*, D0 Collaboration, Phys. Rev. Lett. **105**, 191802 (2010).
- [16] T. Aaltonen *et al.*, CDF Collaboration, arXiv:1105.2815 [hep-ex], (2011). Accepted by Phys. Rev. Lett.